

# ISSUES IN HIGGS PHYSICS

## LECTURE 2

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## SM SATISFIES PERTURBATIVE UNITARITY

- Consider  $2 \rightarrow 2$  particle elastic scattering

$$\frac{d\sigma}{d\Omega} = \frac{1}{64\pi^2 s} |A|^2$$

- Partial wave decomposition of amplitude

$$A = 16\pi \sum_{l=0}^{\infty} (2l+1) P_l(\cos \theta) a_l$$

- $a_l$  are the spin  $l$  partial waves

## UNITARITY

- $P_l(\cos\theta)$  are Legendre polynomials:

$$\int_{-1}^1 dx P_l(x) P_{l'}(x) = \frac{2\delta_{l,l'}}{2l+1}$$

$$\sigma = \frac{8\pi}{S} \sum_{l=0}^{\infty} (2l+1) \sum_{l'=0}^{\infty} (2l'+1) a_l a_{l'}^* \int_{-1}^1 d \cos \theta P_l(\cos \theta) P_{l'}(\cos \theta)$$

$$= \frac{16\pi}{S} \sum_{l=0}^{\infty} (2l+1) |a_l|^2$$

Sum of positive definite terms

## MORE ON UNITARITY

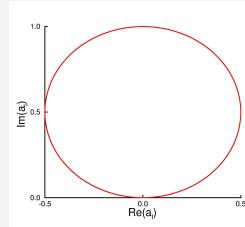
- Optical theorem

$$\sigma = \frac{1}{S} \text{Im}[A(\theta = 0)] = \frac{16\pi}{S} \sum_{l=0}^{\infty} (2l+1) |a_l|^2$$

$$\text{Im}(a_l) = |a_l|^2$$

Optical theorem derived assuming  
only conservation of probability

- Unitarity requirement:

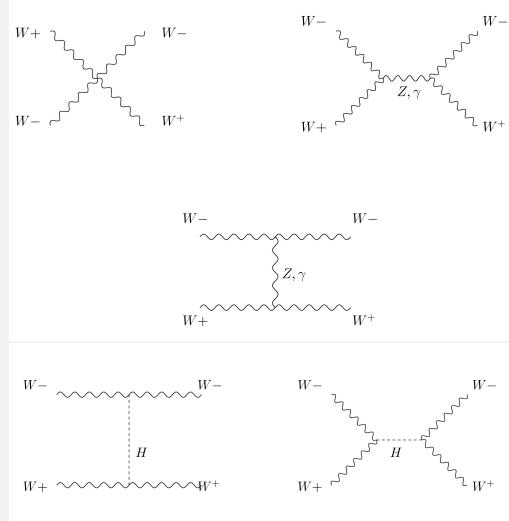


$$|\text{Re}(a_l)| \leq \frac{1}{2}$$

True for any 2→2  
scattering amplitude

Unitarity is major restriction on any model of new physics

## SM IS SPECIAL AT HIGH ENERGY



$E^4$  terms cancel between TGC and QGC

$$A \approx g^2 \frac{E^2}{M_W^2}$$

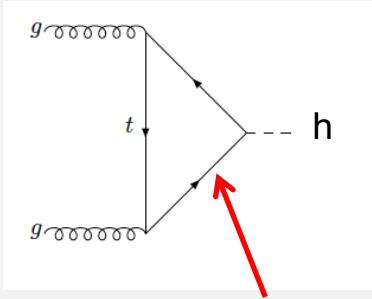
$$A \approx -g^2 \frac{E^2}{M_W^2}$$

Terms which grow with energy cancel for  $E \gg M_h$

SM particles have just the right couplings so amplitudes don't grow with energy

# EXPLORING THE HIGGS SECTOR

# HIGGS PRODUCTION AT A HADRON COLLIDER



Depends on new physics in loop

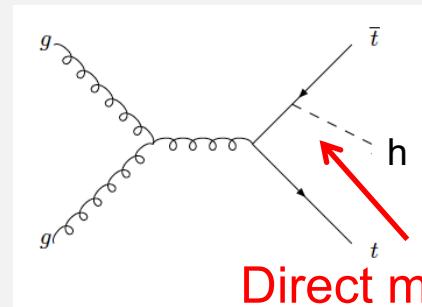
Most important processes:

$$gg \rightarrow h$$

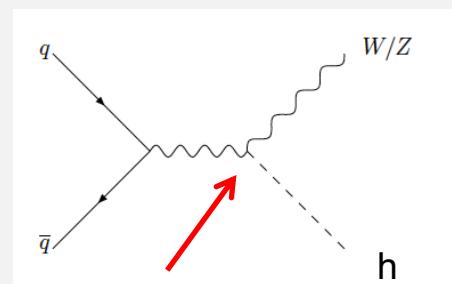
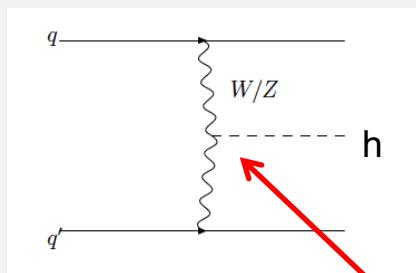
$$q\bar{q} \rightarrow q\bar{q}h$$

$$q\bar{q} \rightarrow q\bar{q}h$$

$$q\bar{q}, gg \rightarrow t\bar{t}h$$



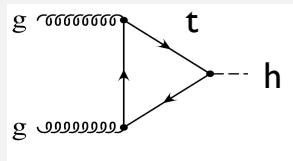
Direct measurement of tth Yukawa



Vanishes if v=0: Fundamental test of EWSB mechanism

## HOW DO WE MAKE A HIGGS BOSON?

- Largest production rate is from gluon fusion
- Largest contribution in SM is from top quarks
- Not a direct measurement of  $t\bar{t}h$  coupling since there could be new particles in loop



$$A(gg \rightarrow h) = -\frac{\alpha_s}{\pi v} \delta_{AB} \left( g^{\mu\nu} \frac{M_h^2}{2} - p^\nu q^\mu \right) \int dx dy \left( \frac{1 - 4xy}{1 - M_h^2 xy/m_t^2} \right)$$

$$L \rightarrow -\frac{\alpha_s}{12\pi v} h G_{\mu\nu}^A G^{\mu\nu,A}$$

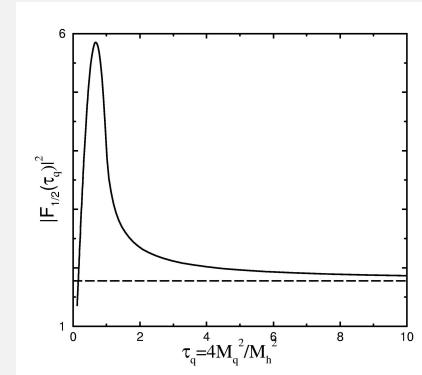
- Independent of  $m_t$  for  $m_t \gg M_h$ : Form useful for radiative corrections

No decoupling of heavy top since coupling proportional to mass

## GLUON FUSION

$$\sigma(gg \rightarrow h)(s) = \frac{\alpha_s^2}{1024\pi v^2} \left| \sum_q F_{1/2}(\tau_q) \right|^2 \delta\left(1 - \frac{M_h^2}{s}\right)$$

- Lowest order cross section:
  - $\tau_q = 4m_q^2/M_h^2$
  - Light Quarks:  $F_{1/2} \rightarrow (m_b/M_h)^2 \log^2(m_b/M_h)$
  - Heavy Quarks:  $F_{1/2} \rightarrow -4/3$



- Rapid approach to heavy quark limit: Counts number of heavy fermions

Heavy quarks don't decouple

# WHAT IF THERE WERE A SM 4<sup>TH</sup> GENERATION?

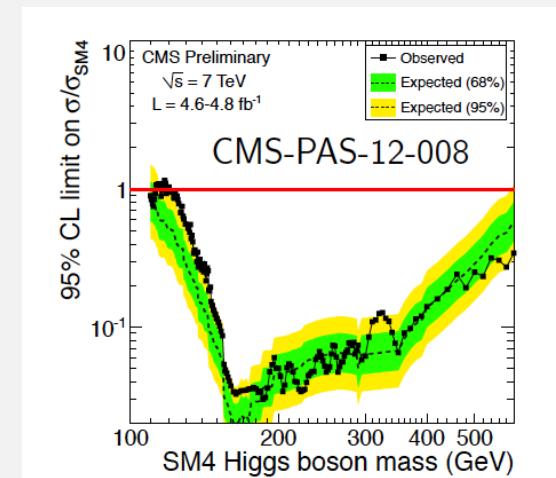
- Each fermion would contribute the same to  $gg \rightarrow h: t, T, B$  (in heavy fermion limit)

$$\sigma \rightarrow \sigma_{SM} (1 + 1 + 1)^2 \rightarrow 9\sigma_{SM}$$

Important: Contribution from chiral fermions roughly independent of fermion mass

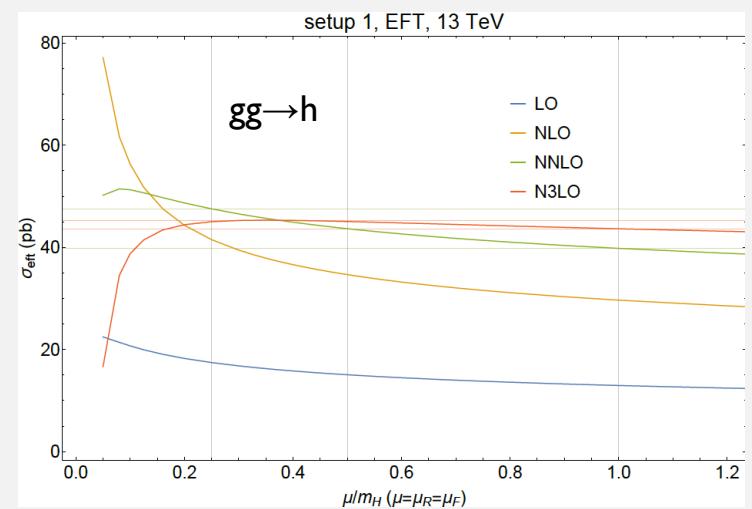
SM 4<sup>th</sup> generation is ruled out by Higgs observation!

\* This logic doesn't hold for vector-like fermions



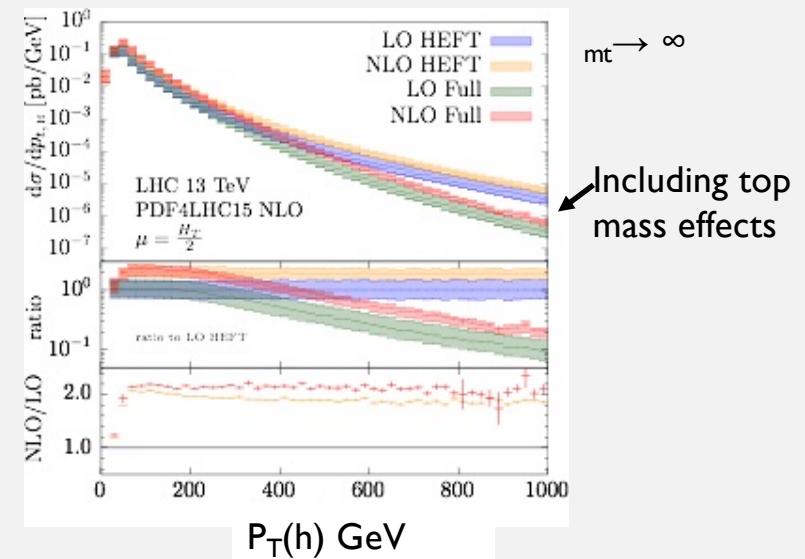
## TREMENDOUS THEORY EFFORT

- MANY (!) Higgs processes calculated to NNLO and some further
- Since Higgs couplings are normalized by the SM expectations, theory and experimental precision must go hand in hand
- These higher order corrections are numerically significant

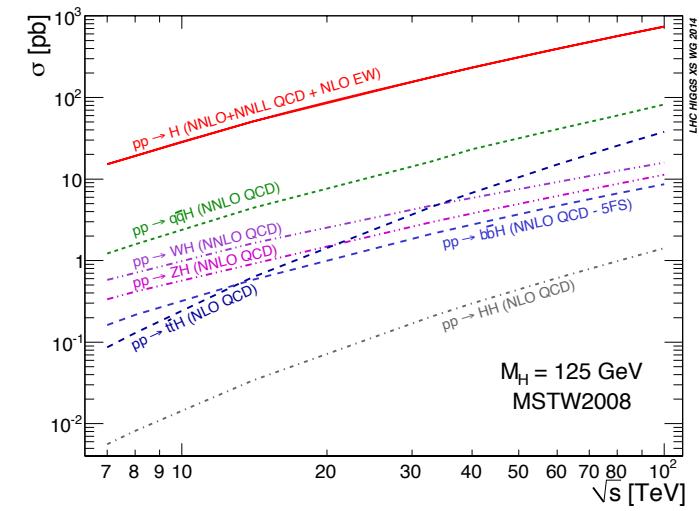


## HIGGS PLUS JET PRODUCTION

- At lowest order Higgs has no  $p_T$
- At high  $p_T$ , can't calculate using  $m_t \rightarrow \infty$
- For  $p_T > m_t$  top quark loop is resolved



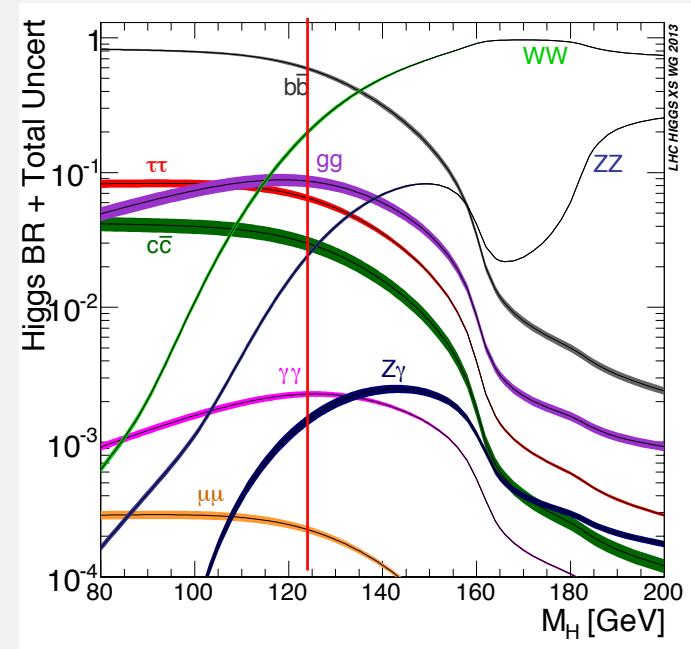
## PRECISE PREDICTIONS FOR PRODUCTION AND DECAY



- Rates to NNLO or NLO
- Gluon fusion dominates
- Rates increase with energy
- $t\bar{t}H$  and  $hh$  smallest rates

## HIGGS DECAYS AND WIDTH

- $\Gamma_{\text{SM}}=4 \text{ MeV} \ll \text{detector resolution}$
- Width is sensitive to light “invisible” particles



## WHAT ARE THE MISSING PIECES?

- Fermion couplings to  $b, t, \tau$
- Gauge boson couplings to  $W/Z/g/\gamma$
- Higgs is spin 0
- No measurements of  $hZ\gamma$ , 2<sup>nd</sup> generation fermions,  $h^3, h^4$  couplings....
- Higgs  $h^2$  coupling (Mass)

What do we want to learn from future Higgs measurements?

## WHAT DO WE LEARN FROM $M_h$ ?

- Free parameter in SM  
(Run-I combination)

$$M_h = 125.09 \pm .21(stat) \pm .11(sys) \text{ GeV}$$

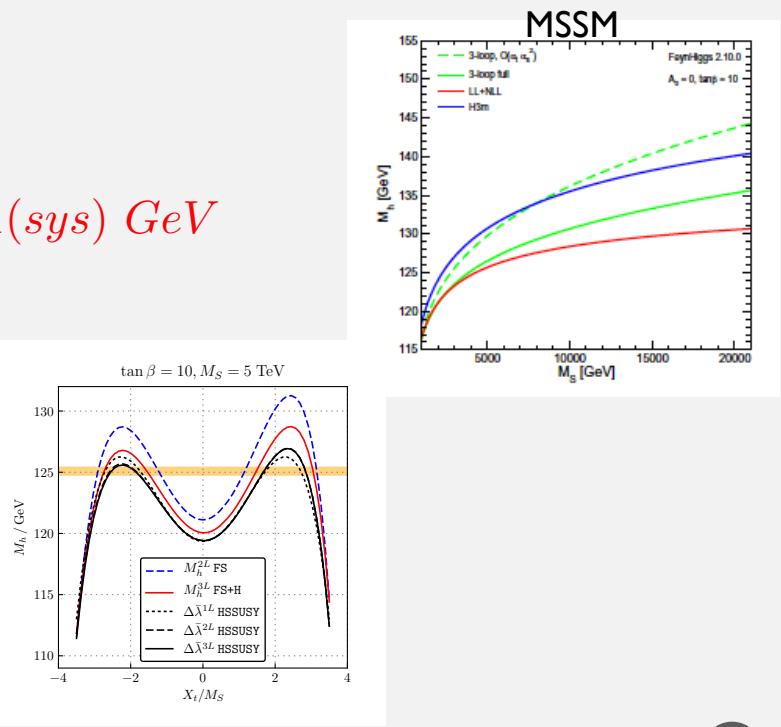
- Higgs self coupling is perturbative:

$$\lambda = \frac{M_h^2}{2v^2} = .13$$

- Consistent with predictions of weak scale MSSM

Hahn, Heinemeyer, Hollik, Rzezak, Weiglein, 2016

Harlander, Klappert, Ochoa, Voigt, 1807.03509



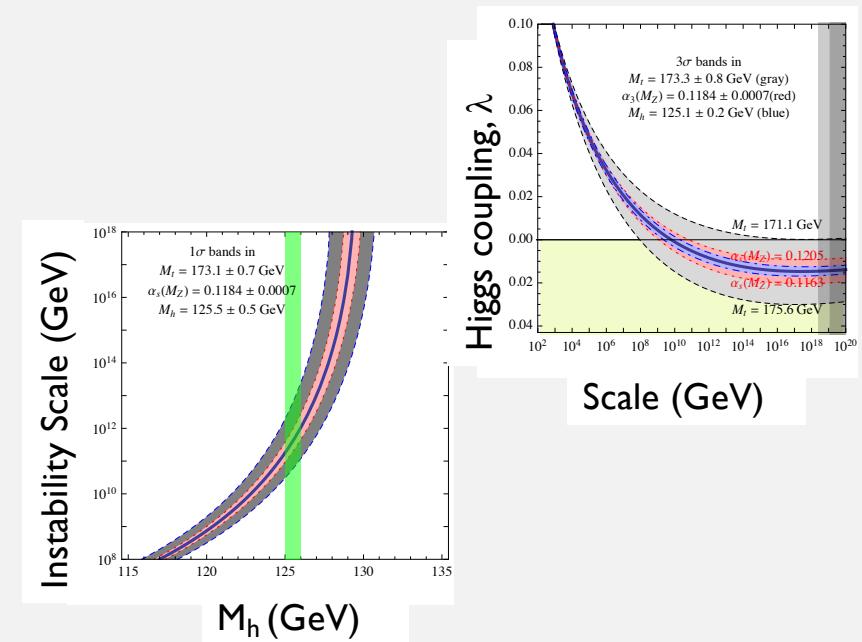
## CAN CALCULATE RUNNING OF HIGGS SELF-COUPLING

- Renormalization group scaling

$$16\pi^2 \frac{d\lambda}{d \log \mu^2} = 12\lambda^2 + 12\lambda y_t^2 - 12y_t^4 + \dots$$

- Large  $\lambda$  (Heavy Higgs):* self coupling causes  $\lambda$  to grow with scale
- Small  $\lambda$  (Light Higgs):* coupling to top quark causes  $\lambda$  to become negative

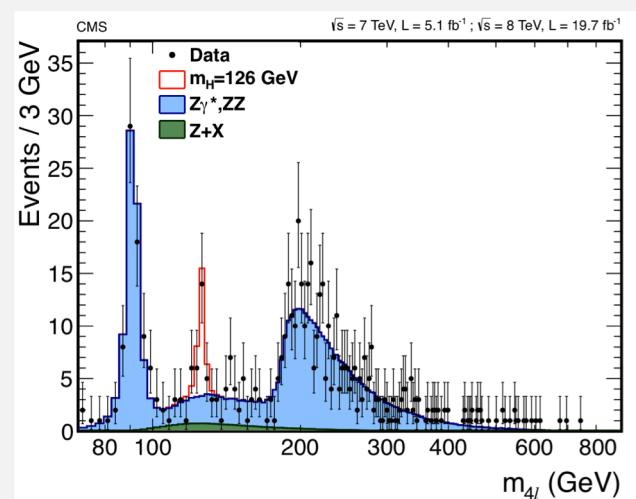
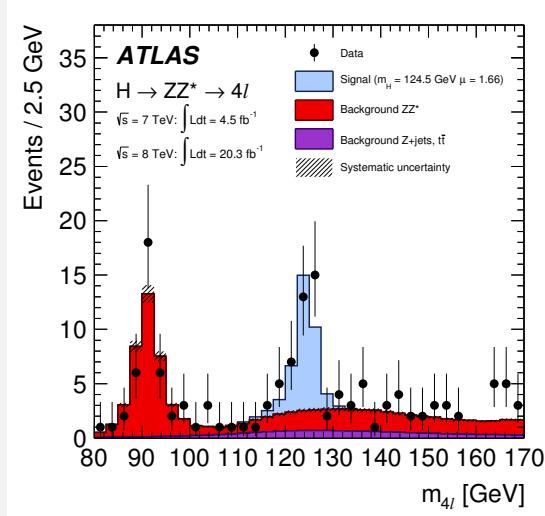
$$y_t = \frac{M_t}{v}$$



## WHAT DOES THE HIGGS WIDTH TELL US?

$gg \rightarrow h \rightarrow ZZ$

- Goal: Measure  $gg \rightarrow h \rightarrow ZZ$  and use insights about resonances

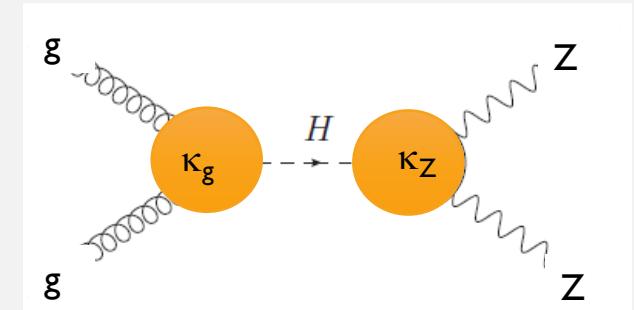


## HIGGS RESONANCE

$$\sigma \sim \int ds \frac{\kappa_g^2 \kappa_Z^2}{(s - M_h^2)^2 + M_h^2 \Gamma_h^2} \epsilon_Z^\mu(p) \cdot \epsilon_Z^{\mu*} \epsilon_Z^\nu(q) \cdot \epsilon_Z^{\nu*}(q)$$

Longitudinal Z polarization grows with energy:  $\epsilon_L^\mu \sim \frac{p^\mu}{M_Z}$

Consider above pole,  $s \gg M_h^2$        $\sigma_{above} \sim \frac{\kappa_g^2 \kappa_Z^2}{M_h^2}$



No dependence on width

## HIGGS RESONANCE

$$\sigma \sim \int ds \frac{\kappa_g^2 \kappa_Z^2}{(s - M_h^2)^2 + M_h^2 \Gamma_h^2} \epsilon_Z^\mu(p) \cdot \epsilon_Z^{\mu*} \epsilon_Z^\nu(q) \cdot \epsilon_Z^{\nu*}(q)$$

$$\int ds \frac{1}{(s - M_h^2)^2 + M_h^2 \Gamma_h^2} \rightarrow \frac{\pi}{\Gamma_h M_h} \delta(s - M_h^2)$$

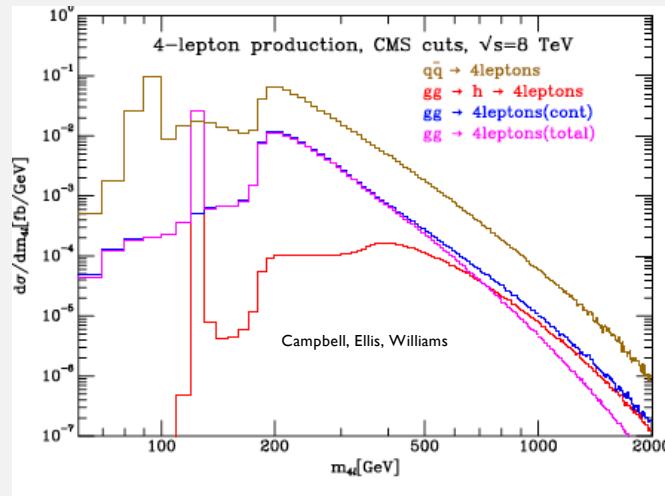
On peak:  $\sigma_{res} \sim \frac{\kappa_g^2 \kappa_Z^2}{\Gamma_h M_h}$

Resonance rate sensitive  
to total Higgs width

IDEA



- Measure above and below the peak:



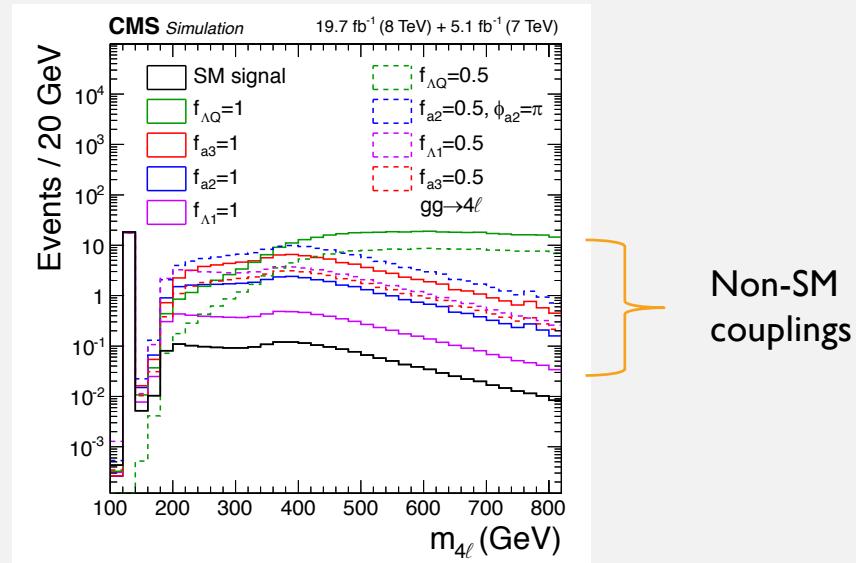
$$\frac{\sigma_{above}}{\sigma_{res}} \sim \Gamma_h$$

About 15% of total cross section in  $m_{4l} > 140$  GeV region above peak

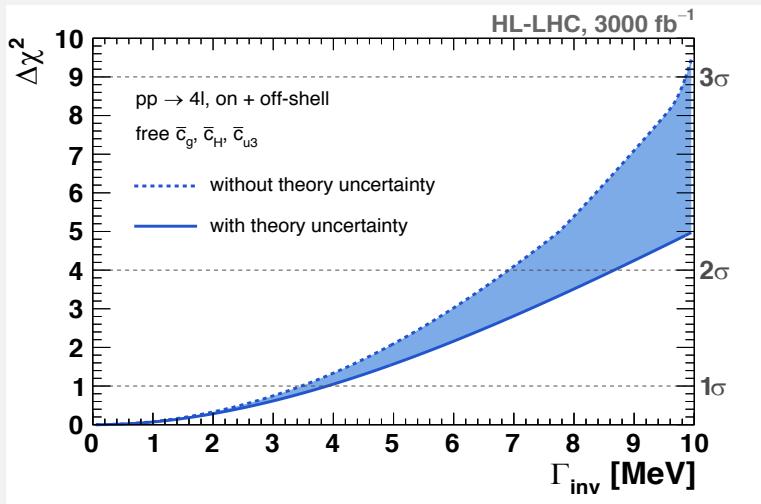
ATLAS, CMS,  $\Gamma_h < (4 - 5)\Gamma_{SM}$

## IMPLICIT ASSUMPTIONS IN HIGGS WIDTH EXTRACTION

- Assumes couplings are same on resonance and off resonance
- Assumes SM, ie no new higher dimension operators



## INVISIBLE HIGGS WIDTH



Interpret Higgs width limit as limit on invisible width, allowing non-SM couplings to vary

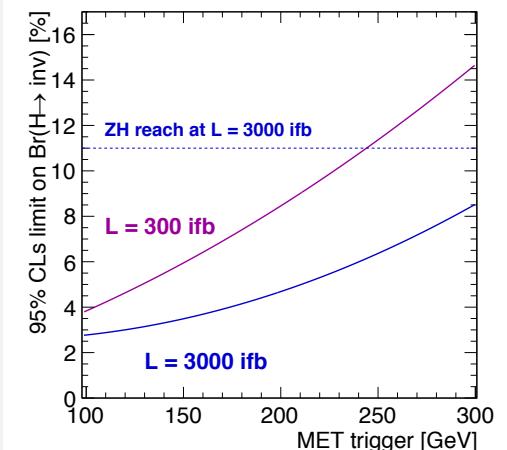
SM invisible width is  $\text{BR } (h \rightarrow vvvv) = 10^{-3}$

## INVISIBLE WIDTH

- Sources of invisible width:
  - Light scalars, light axions, relaxions....  $h \rightarrow \phi_i \phi_i$
  - Neutralinos, fermionic dark matter...  $h \rightarrow \chi \bar{\chi}$
- Limit comes from boosted Z recoiling against large missing energy

- Experimental limit is primarily from  $Z h$ ,  $Z \rightarrow \text{leptons}$
- Future limits from VBF
- $\text{BR}(h \rightarrow \text{invisible}) < 24\%$

Future limits on  
invisible width from VBF



Missing energy  
trigger in VBF

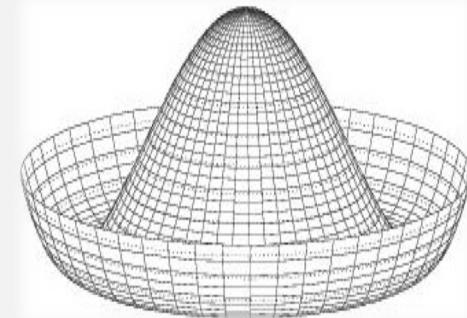
## HIGGS SELF-COUPLING BIG MILESTONE

- We don't know that the Higgs comes from the scalar potential

$$\begin{aligned}V &= \mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2 \\&= \frac{M_h^2}{2} h^2 + \lambda_3 h^3 + \lambda_4 h^4\end{aligned}$$

- SM is perturbative

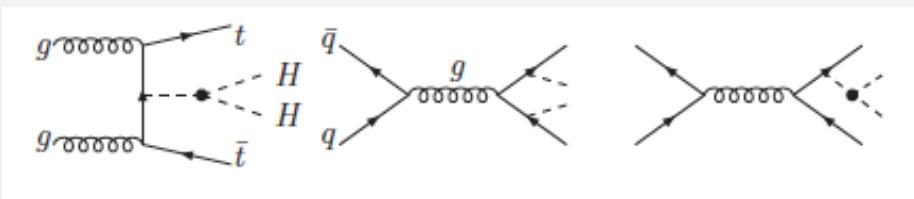
$$\lambda_3 = \frac{M_h^2}{2v} \sim .13v, \quad \lambda_4 = \frac{M_h^2}{8v^2} = .03$$



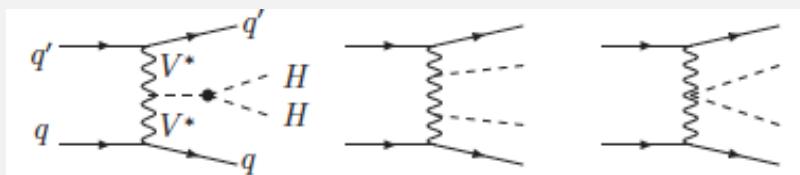
## PRODUCTION OF $hh$



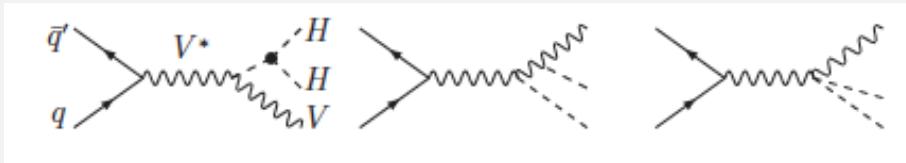
Sensitive to heavy colored particles  
(eg stops or top partners)



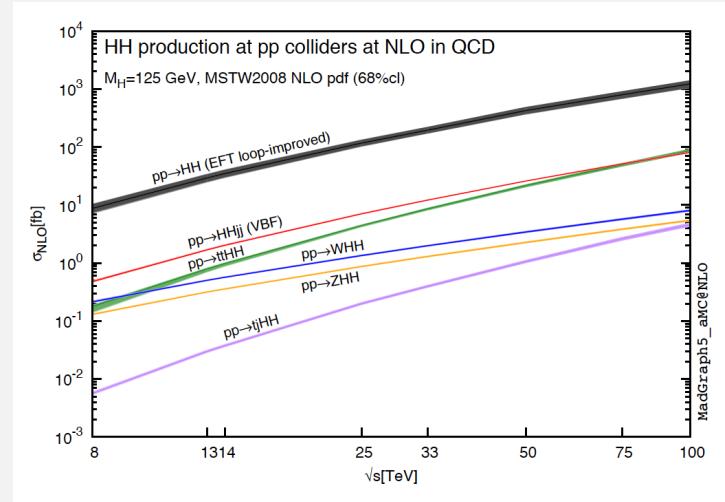
Sensitive to anomalous  
top-Higgs couplings



Sensitive to anomalous VVhh  
couplings



## SMALL RATES FOR $hh$ PRODUCTION



*This is  $3000 \text{ fb}^{-1}$  physics and motivation for 100 TeV Collider!*

## MEASURING THE $hh$ COUPLING

- $hh$  first occurs at one-loop

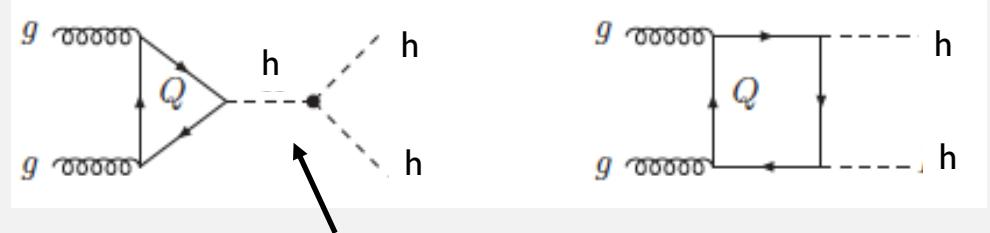
Large cancellation between diagrams

$3 \text{ ab}^{-1}$

Theorists:  $0.7 < \lambda_3/\lambda_{3,\text{SM}} < 7.7$

ATLAS:  $0.2 < \lambda_3/\lambda_{3,\text{SM}} < 7.$  (bbbb)

ATLAS:  $-0.8 < \lambda_3/\lambda_{3,\text{SM}} < 7.7$  (bb $\gamma\gamma$ )



Goal: Measure  $\lambda_3$

CURRENT

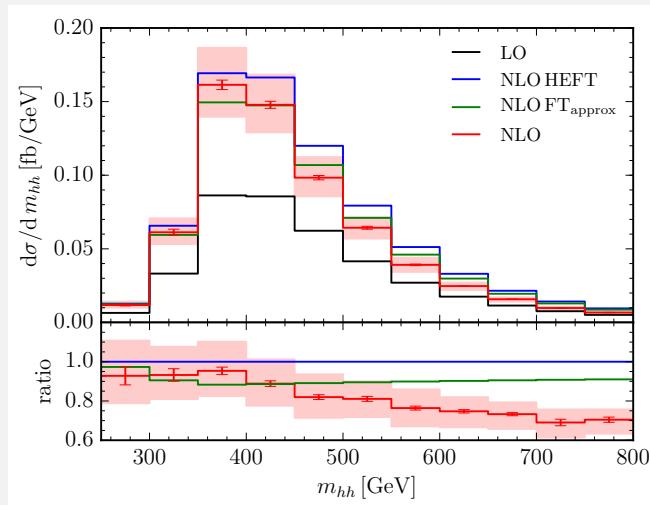
CMS:  $\sigma/\sigma_{\text{sm}} < 19$  (bb $\gamma\gamma$ )

ATLAS:  $\sigma/\sigma_{\text{sm}} < 13.0$  (bbbb)

CMS:  $\sigma/\sigma_{\text{sm}} < 28$  (bb $\tau\tau$ )

## PROGRESS IN HH CALCULATIONS

$gg \rightarrow hh @ 14\text{ TeV}$



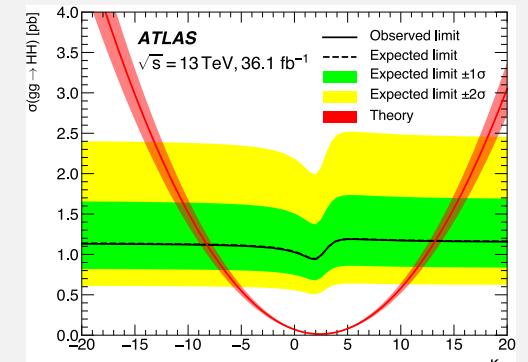
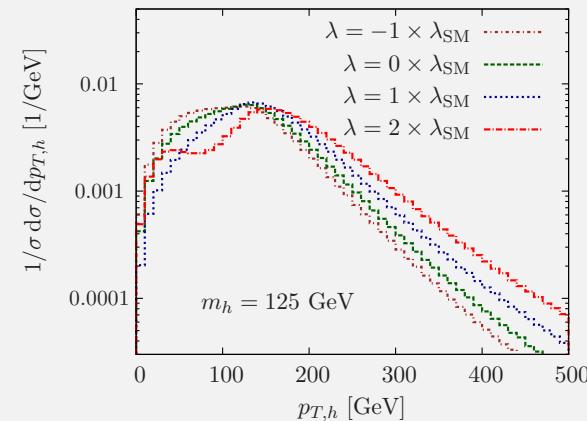
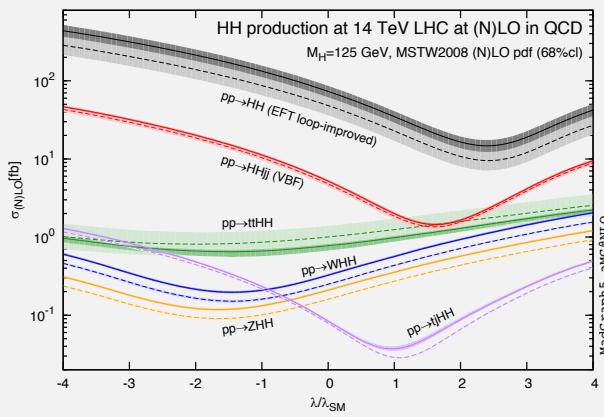
Result with full mass dependence  
20-30% below HEFT ( $m_t \rightarrow \infty$ )  
for  $m_{hh}$  above 450 GeV

LO and  $m_t \rightarrow \infty$  not accurate  
for tails of distributions

→ Need higher order corrections

## MEASURING $\lambda_3$

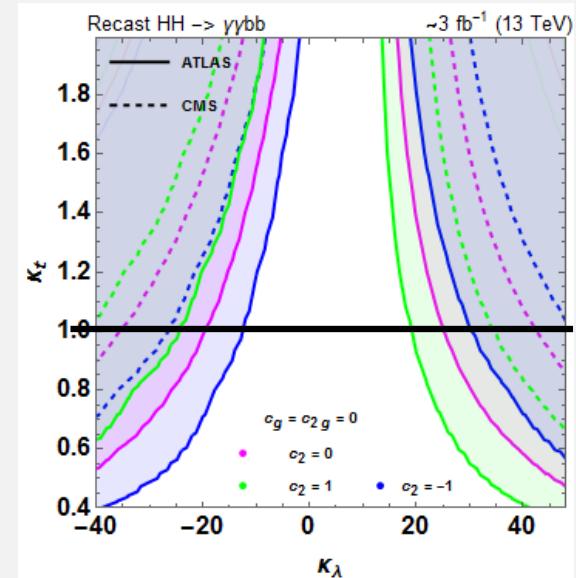
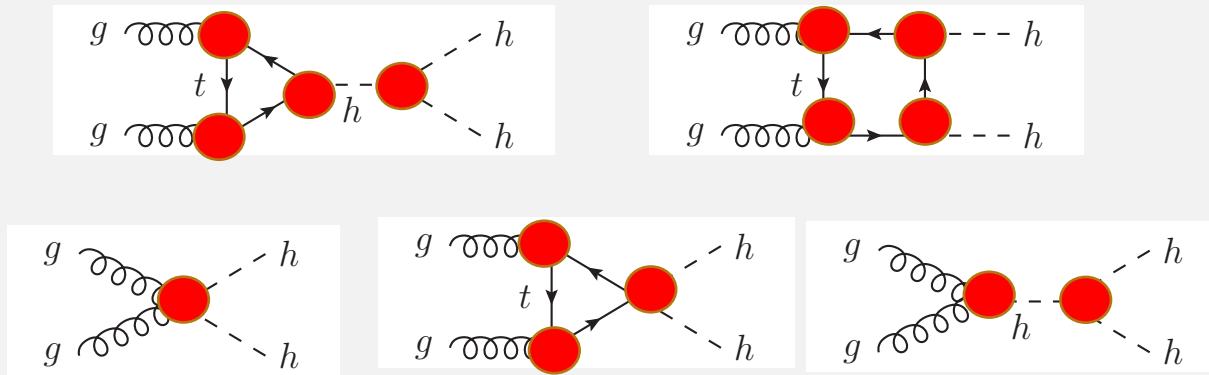
- Corrections to  $\lambda_3$  are typically  $O(v^2/\Lambda^2)$ , ie  $\sim 5\%$  (for  $\Lambda = 1 \text{ TeV}$ ) **(not factors of 2-3)**
- Unitarity from  $hh \rightarrow hh$ :  $\lambda_3/\lambda_3(\text{SM}) < 6.5$



$$\lambda_3/\lambda_3(\text{SM}) = \kappa_\lambda$$

BUT SEE EFT SECTION.....

- In any realistic model, other interactions than  $hhh$  are also changed
- Need global fit: Limits are correlated



## MEASURING THE HIGGS SELF COUPLING INDIRECTLY

- Higgs self coupling contributes to  $gg \rightarrow h$ ,  $h \rightarrow \gamma\gamma$ ,  $t\bar{t}h$  production
- Use measured single Higgs rates to limit:  $-9.4 < \lambda_3/\lambda_3(\text{SM}) < 17$
- Similar order of magnitude as limit from  $hh$  production
- Include single Higgs production contribution to  $\lambda_3$  in global fit
- (We don't yet have all the theory pieces to do this consistently  
... need full 2 loop electroweak corrections with all possible  
couplings)

